

# A Theory for Strong, Long-Lived Squall Lines

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# OUTLINE

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- 1 Introduction
- 2 Observations
- 3 Squall-line simulations
- 4 Interpretation
- 5 Summary

# 1

## INTRODUCTION

## FEATURES OF SQUALL LINE

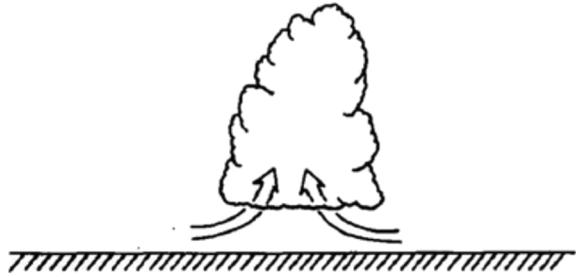
- line or narrow band of active thunderstorms (Glossary of Meteorology)
- lasting several hours

## TWO DISTINCT PREMISES

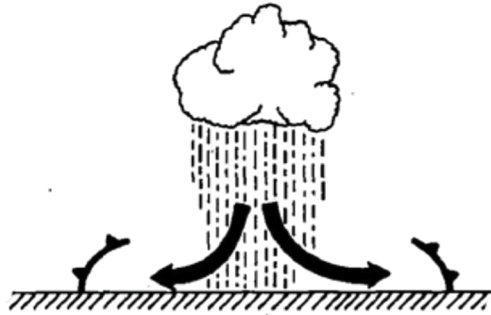
- prone to having a **steady** structure
- occurring in concert along a line (Moncrieff, 1978)

# Introduction

(a)



(b)

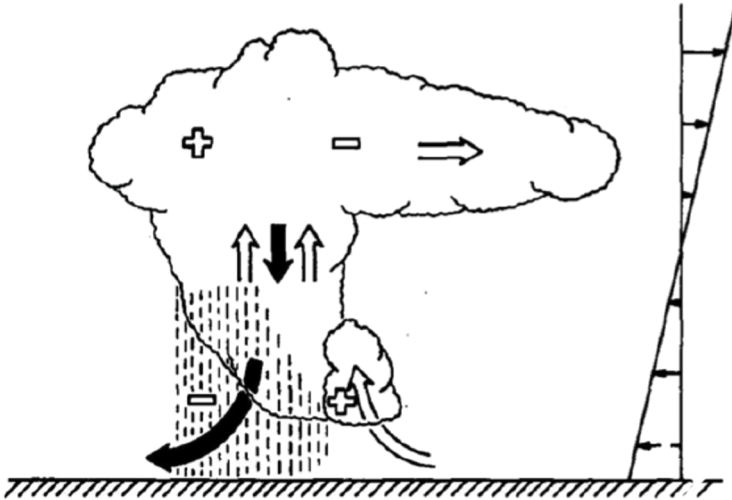


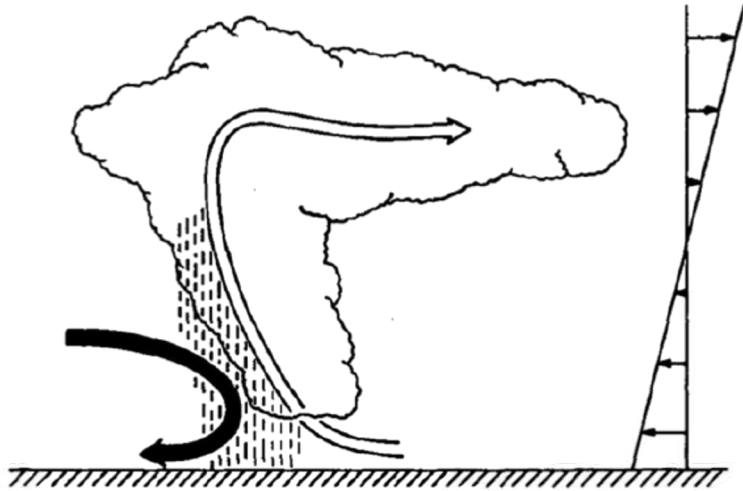
## Byers and Braham(1949)

- A thunderstorm cell has a **life cycle** over which the cell updraft will yield to a downdraft induced by the accumulation of rain within the updraft and, for this reason, a thunderstorm cell is naturally **short-lived** (roughly 30-60 min) .

## Newton(1950)

- squall-line, a system of updrafts and downdrafts aligned perpendicularly to the shear
- **reducing the shear** within the system
- setting up the **convergence** and **divergence**:
  - produces new cells on the downshear side
  - suppress old cells on the upshear side





## Ludlam-Newton Model

- strong updraft **canting** against the wind shear:
  - allowing the updraft to unload its rain upshear
  - permitting the circulations to continue indefinitely
- squall line, a collection of such **long-lived** cells

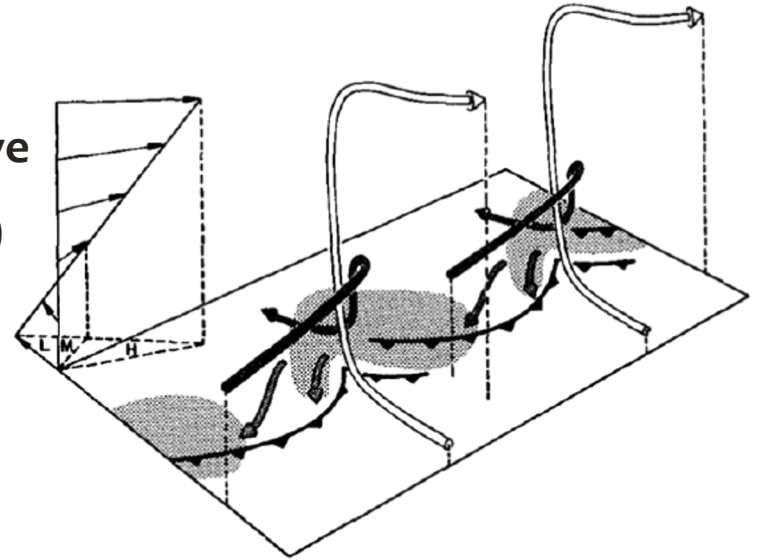
- Modelers **failed** to replicate the Ludlam-Newton observational model in **two-dimensional** simulations.  
due to the inherent **three-dimensionality** of cumulonimbus clouds (Lilly,1979)
- the revised version of the Ludlam-Newton model:  
squall line, a collection of **supercells?**



# Introduction

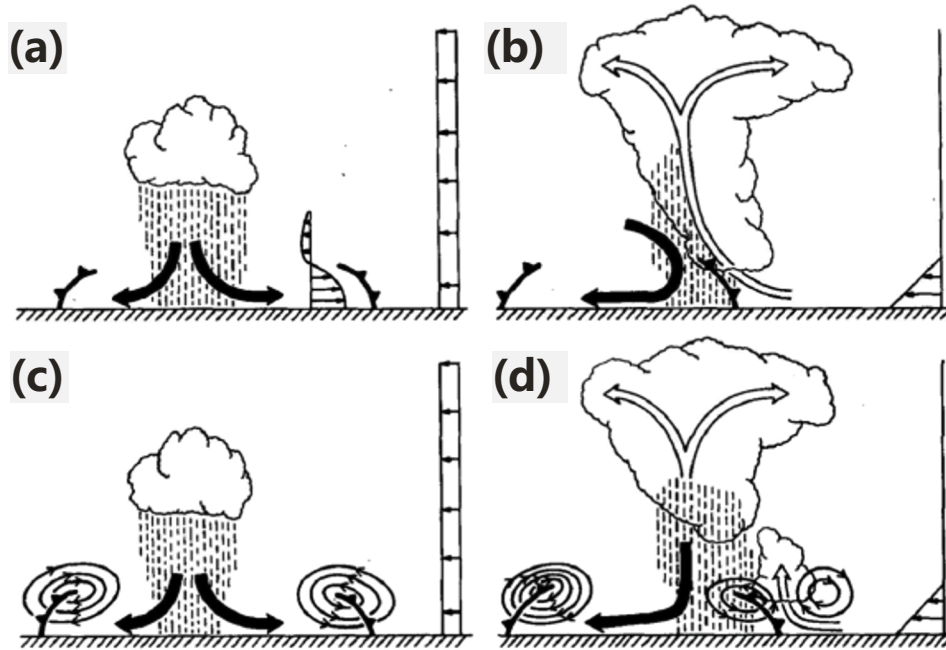
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- A **steady** squall line if they aligned **at an angle to the shear** so that their respective circulations did not interfere. (Lilly,1979)
- Contained **supercell-like** circulations allowed their system to be **long lived**. (Moncrieff and Miller ,1976)
- **However, most squall lines are not composed of supercell thunderstorms.** (Bluestein and Jain, 1985)



# Introduction

09



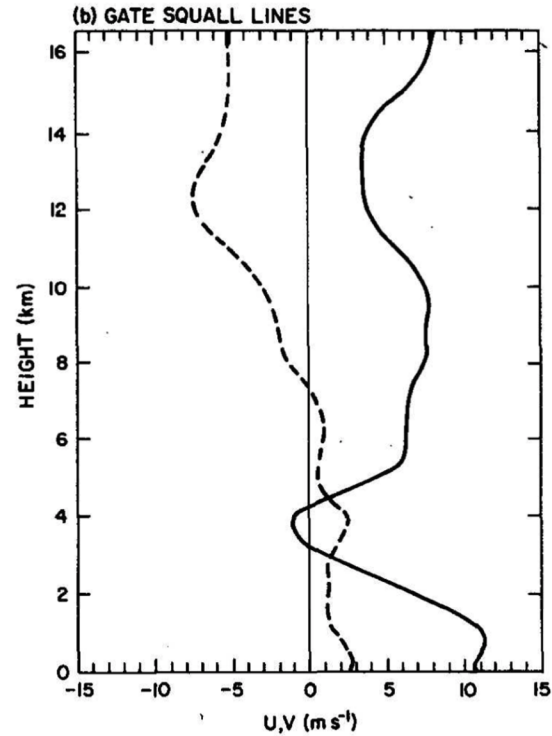
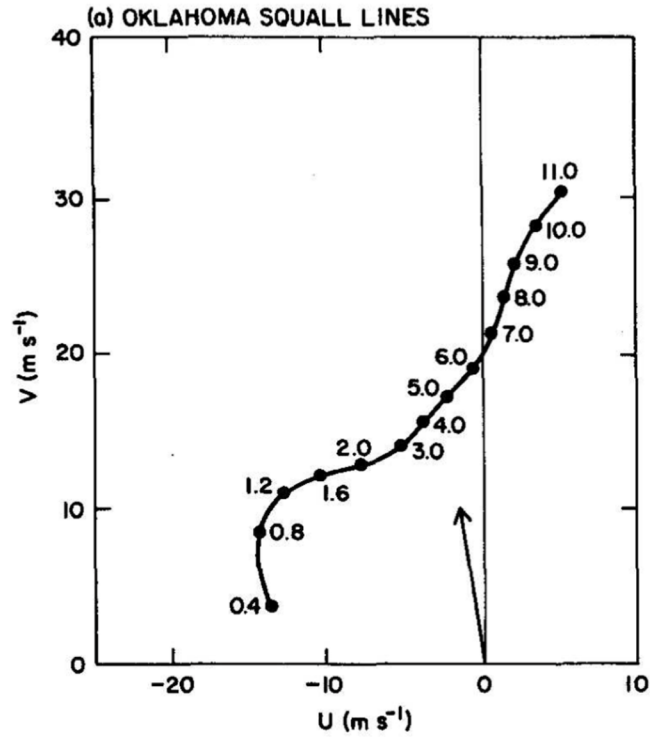
- The long life of the strictly two-dimensional solution is a consequence of the **low-level shear** in the ambient wind profile. (Thorpe et al., 1982; hereafter TMM)

# 2

## OBSERVATIONS

# Observations

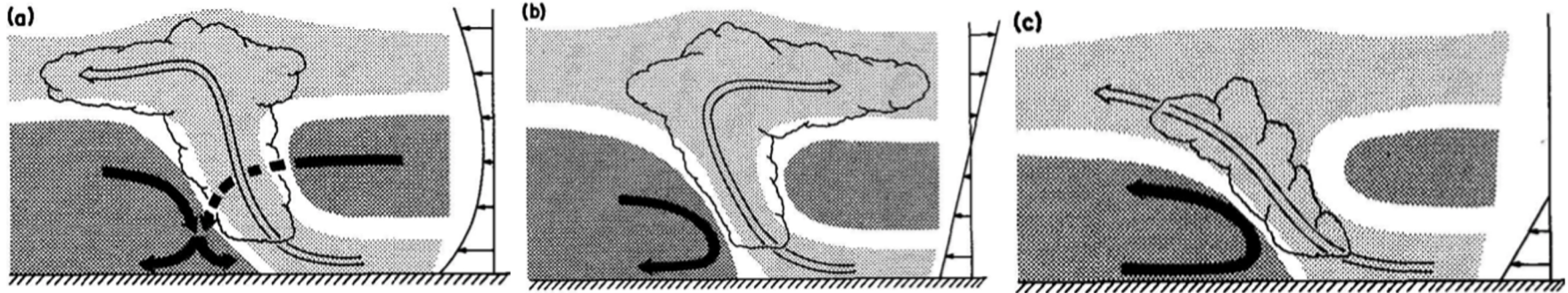
## a. Squall-line Environment



# Observations

## b. Squall-line Circulation

12



### Zipser's analysis

- flow from ahead of to behind the line in the mature phase
- low-level outflow directed towards the front

### Newton's analysis

- all the air approaching the line ascend
- not consider the possibility of substantial three-dimensional motion

### Carbone's analysis

- a "gravity current"
- a wintertime squall line

# 3

## SQUALL-LINE SIMULATIONS

# Squall-line Simulation

## a. Shallow shear normal to the line: two-dimensional simulation

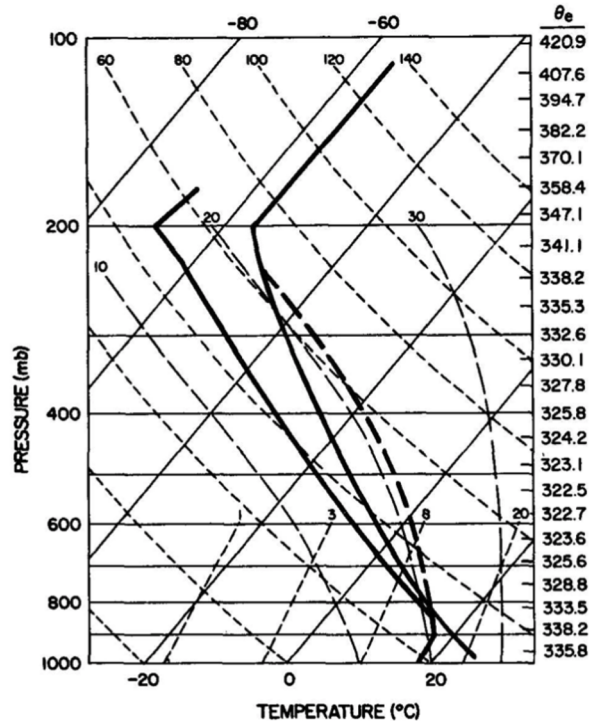
Klemp-Wilhelmson cloud model (hereinafter KW cloud model)

Model setting:

- 180km in x direction; grid interval  $\Delta x = 2\text{km}$
- 17.5km in z direction; grid interval  $\Delta z = 700\text{m}$
- x: open boundary condition
- z: zero-flux-type at low surface
- no ice processes included

# Squall-line Simulation

## a. Shallow shear normal to the line: two-dimensional simulation

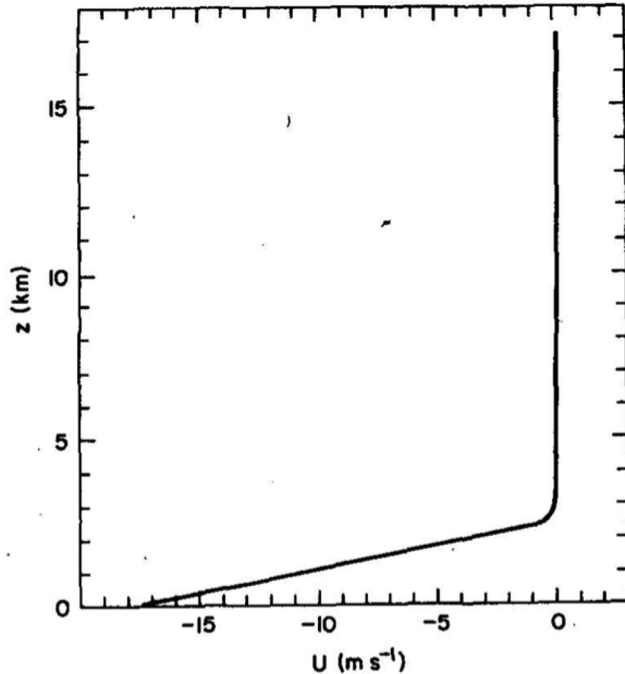


- broadly typical of the environment of midlatitude squall lines
- motion initiated with a warm (2K exceed) line centered at  $x=90$ ,  $z=1.4$  km
- decay to 0 for  $|x-90| > 10$ ;  $|z-1.4| > 1.4$



# Squall-line Simulation

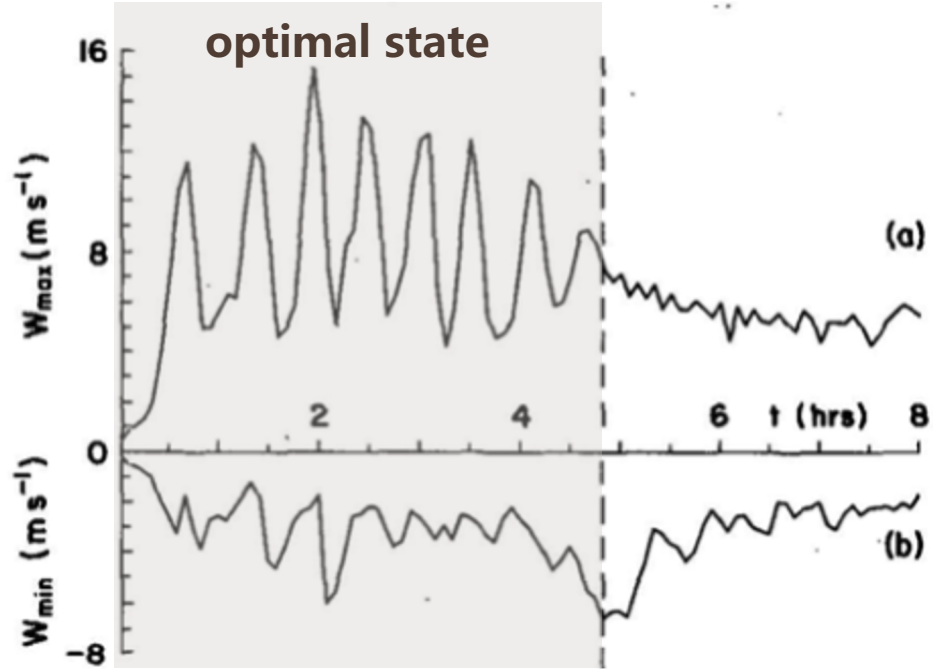
## a. Shallow shear normal to the line: two-dimensional simulation



- A shear of **17.5 m/s** in the lowest 2.5 km allows for the longest sequence of the strongest cell updrafts.

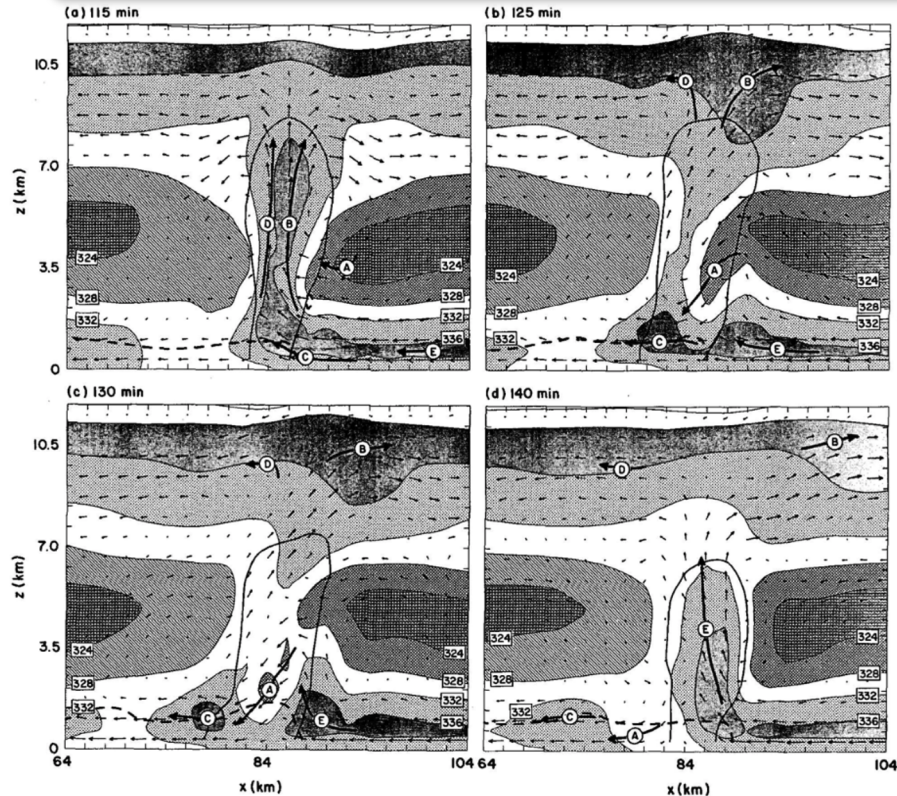
# Squall-line Simulation

## a. Shallow shear normal to the line: two-dimensional simulation



# Squall-line Simulation

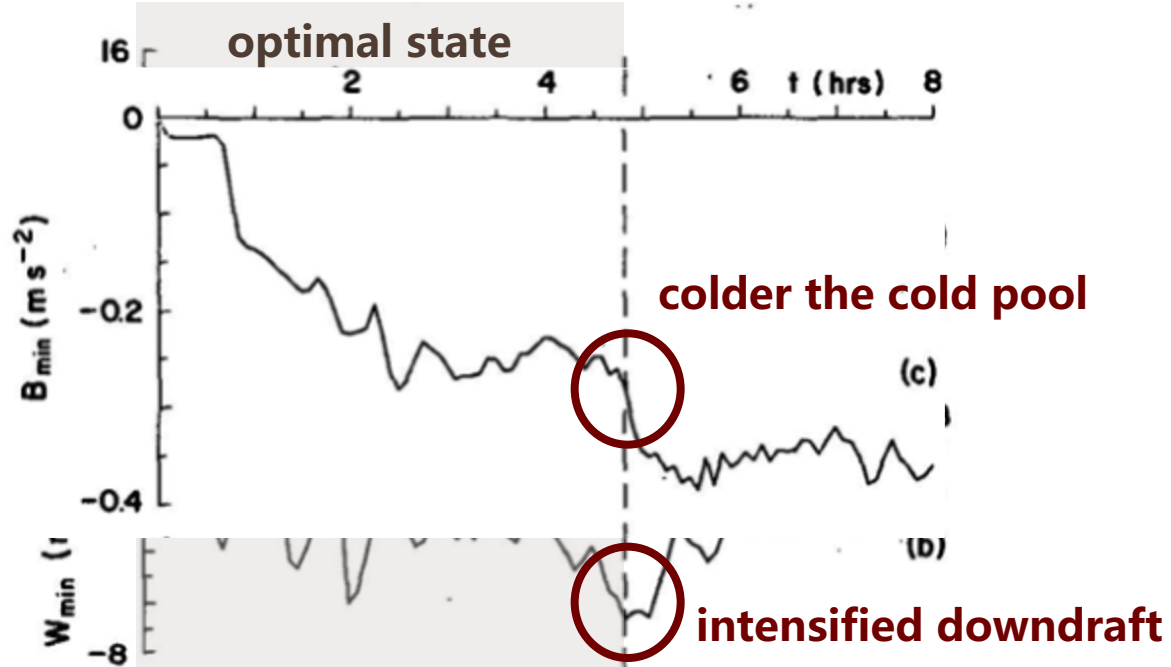
## a. Shallow shear normal to the line: two-dimensional simulation



- optimal state
- thick dashed line :  $-1K$  perturbation  $\theta_e$  contour
- solid line: 2 g/kg rainwater contour
- shaded:  $\theta_e$  field
- grid interval: 16m/s

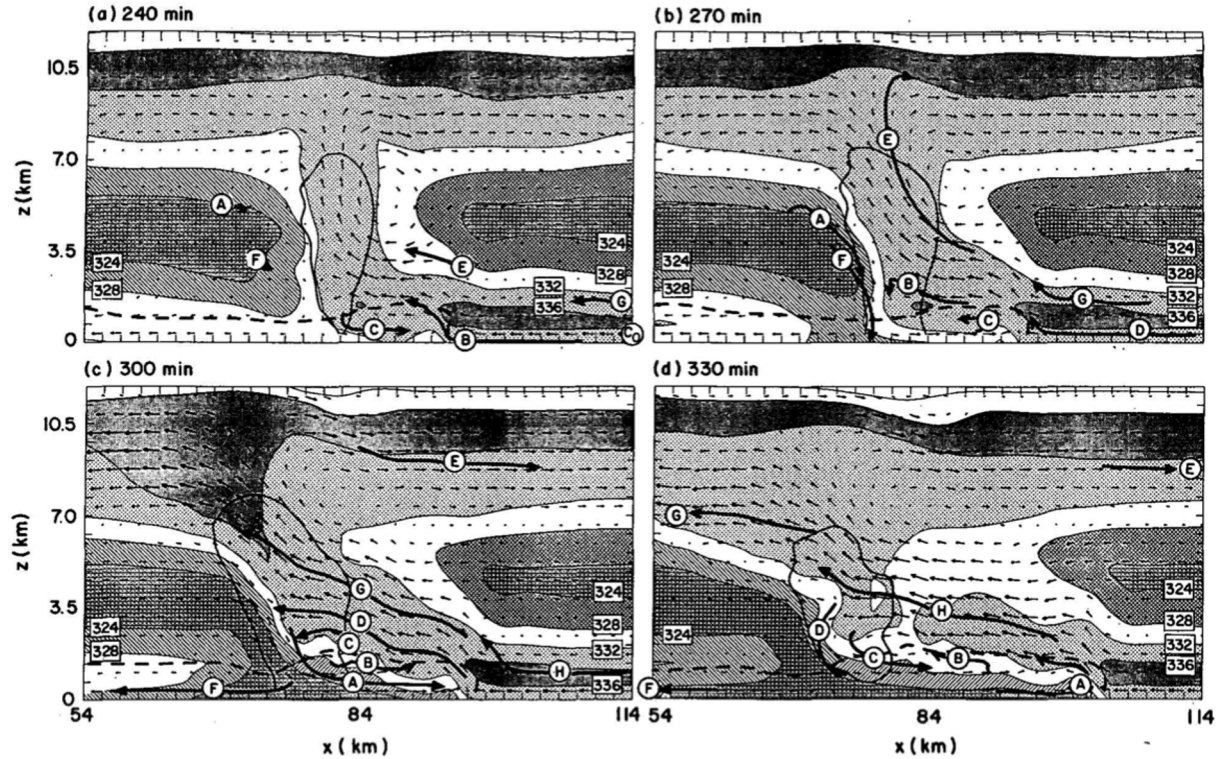
# Squall-line Simulation

## a. Shallow shear normal to the line: two-dimensional simulation



# Squall-line Simulation

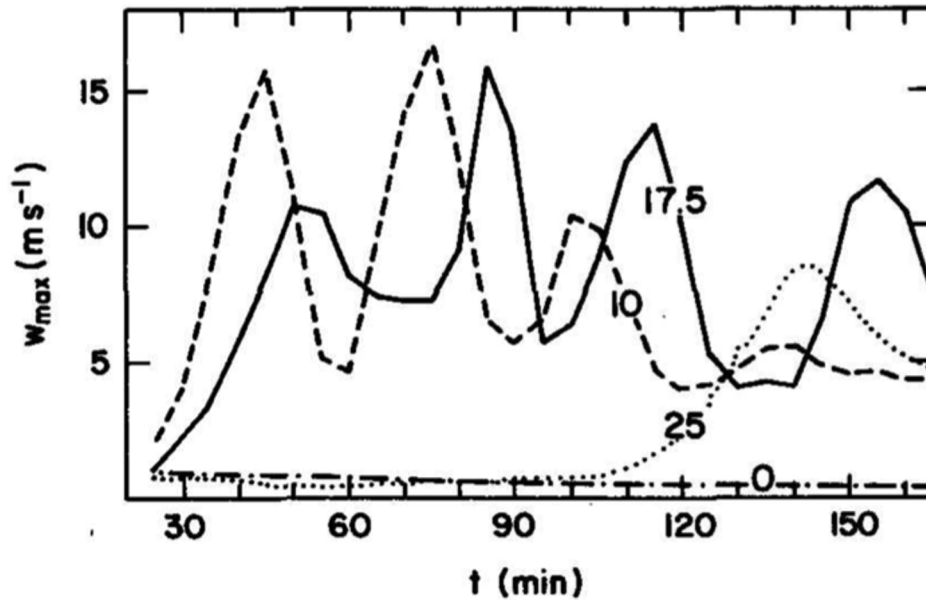
## a. Shallow shear normal to the line: two-dimensional simulation



- weaker mode

# Squall-line Simulation

## a. Shallow shear normal to the line: two-dimensional simulation



- cold-pool initialization

# Squall-line Simulation

## b. Shallow shear normal to the line: three-dimensional simulation

KW cloud model

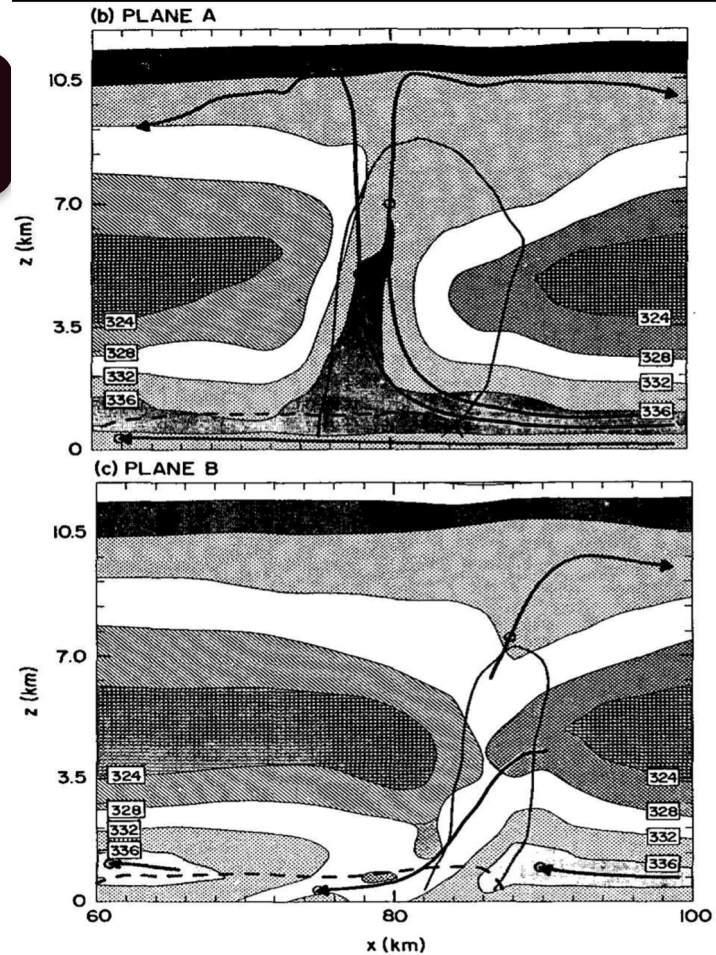
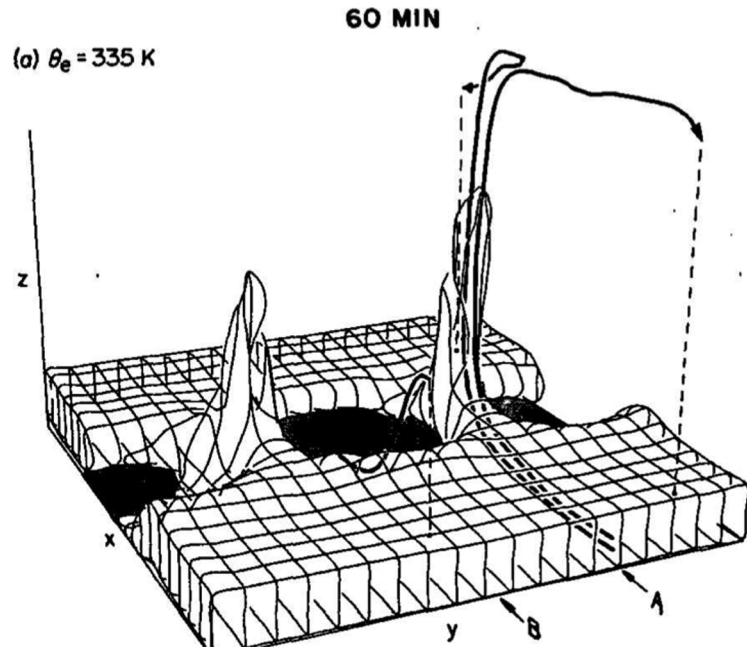
Model setting:

- 120km in y direction; grid interval  $\Delta y = 2\text{km}$
- y: periodic boundary condition
- small ( $< 0.1\text{ K}$ ) random temperature perturbations on the initiating line thermal

Other conditions are the same as in the two-dimensional simulation.

# Squall-line Simulation

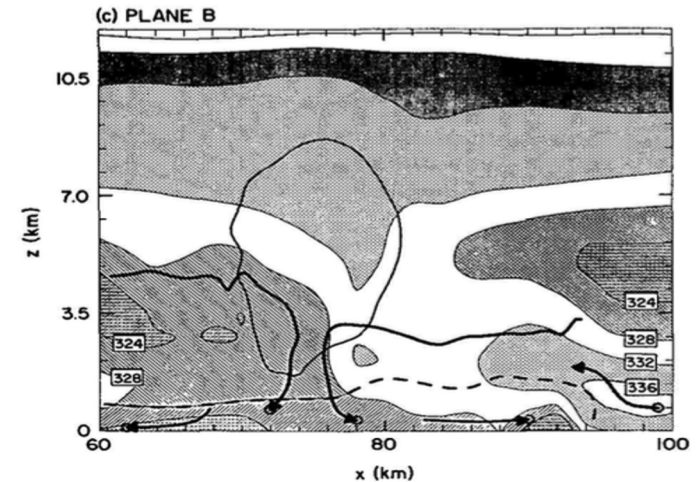
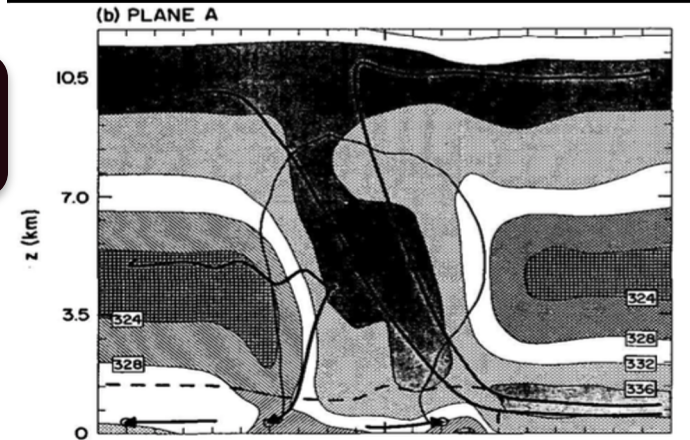
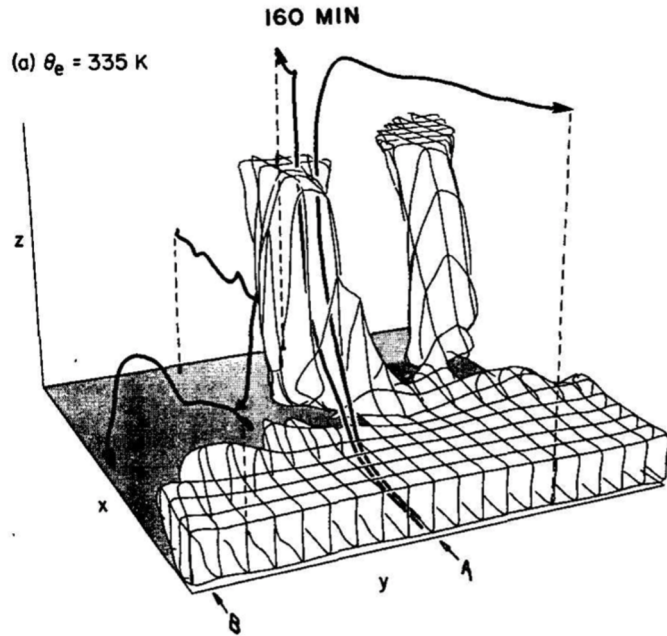
b. Shallow shear normal to the line:  
three-dimensional simulation





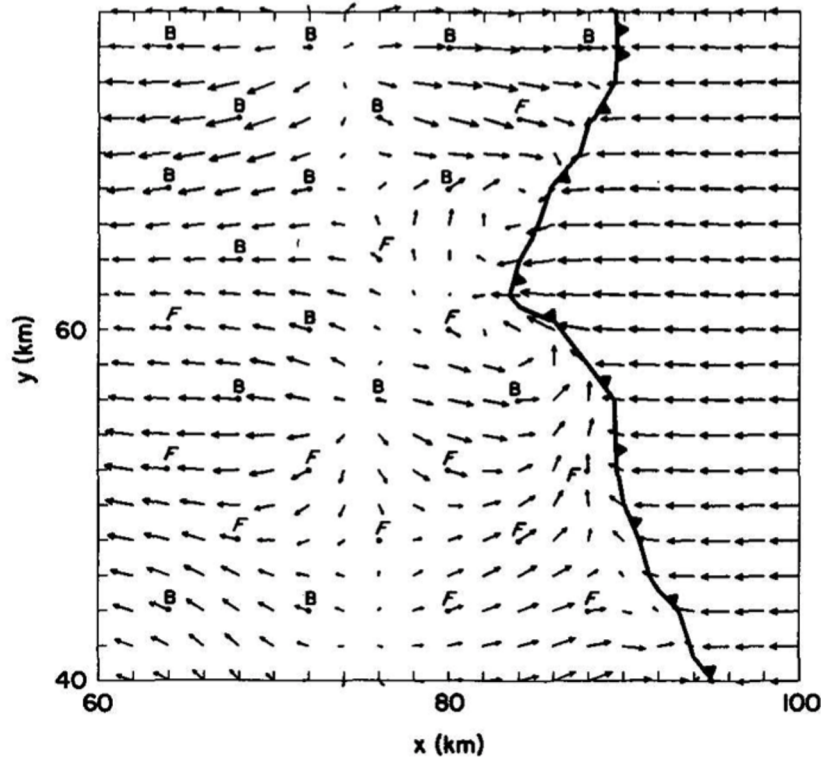
# Squall-line Simulation

b. Shallow shear normal to the line:  
three-dimensional simulation



# Squall-line Simulation

## b. Shallow shear normal to the line: three-dimensional simulation

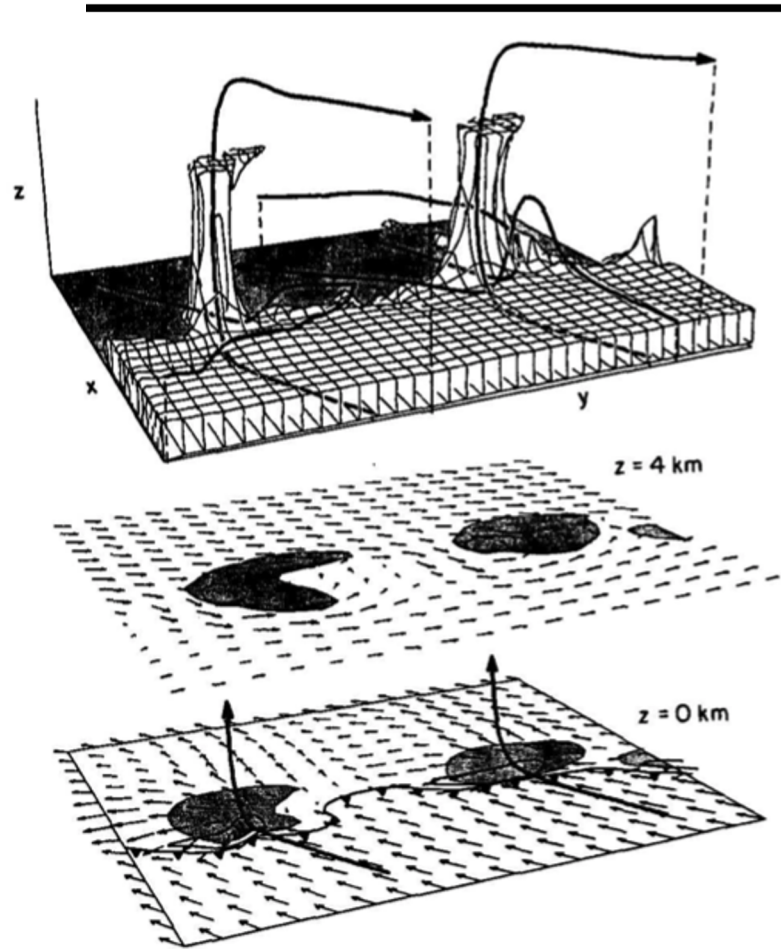


- $z = 350$  km
  - grid interval: 16m/s
  - Barbed line: -1K perturbation
- $\theta_e$  contour

# Squall-line Simulation

## c. Deep shear at $45^\circ$ to the line

- The time-dependence found in the three-dimensional, nonsupercellular squall-line simulations **is not an inevitable feature** of the present numerical simulation.



# 4

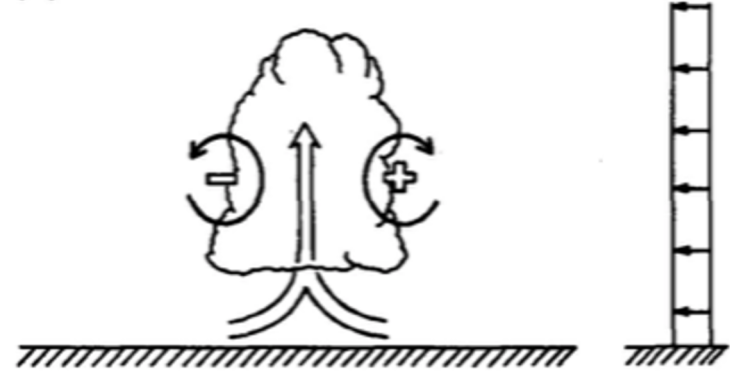
## INTERPRETATION

# Interpretation

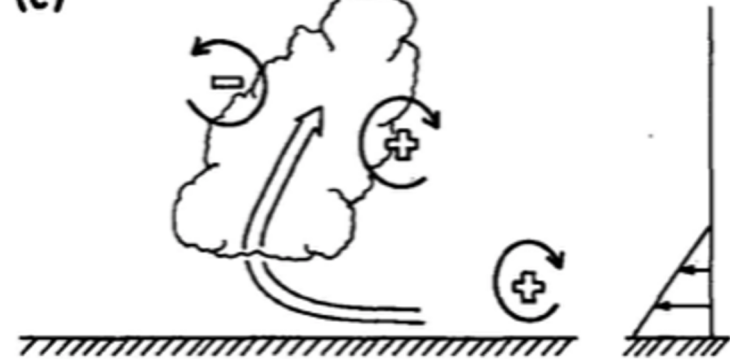
## a. The role of wind shear

$$\frac{d\eta}{dt} = -\frac{\partial B}{\partial x}, \quad \eta = \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x}$$

(a)



(c)

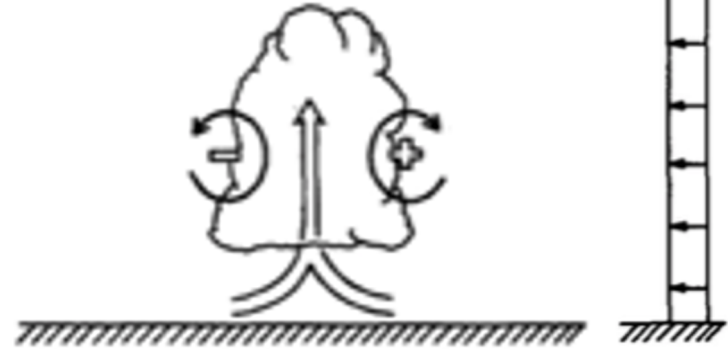


# Interpretation

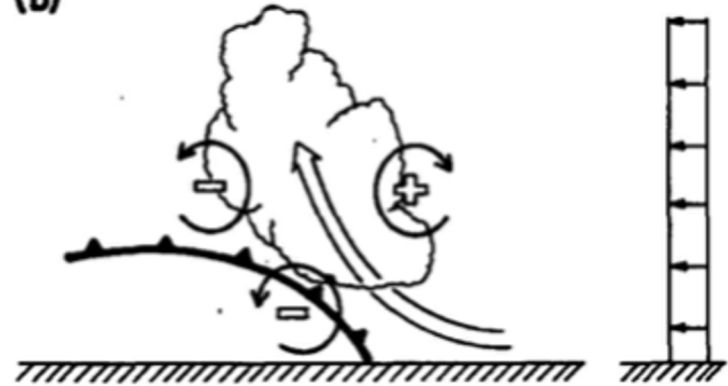
## b. The role of the cold pool

$$\frac{d\eta}{dt} = -\frac{\partial B}{\partial x}, \quad \eta = \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x}$$

(a)



(b)



## c. The cold pool with and without low-level shear

$$\frac{d\eta}{dt} = -\frac{\partial B}{\partial x}, \quad \eta = \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x}$$

$$\text{, with } \eta \frac{\partial u}{\partial x} + \eta \frac{\partial w}{\partial z} = 0$$

$$\frac{\partial \eta}{\partial t} = -\frac{\partial(u\eta)}{\partial x} - \frac{\partial(w\eta)}{\partial z} - \frac{\partial B}{\partial x}$$

$$\begin{aligned} \frac{\partial}{\partial t} \int_0^d \int_L^R \eta \, dx \, dz &= \int_0^d (u\eta)_L \, dz - \int_0^d (u\eta)_R \, dz \\ &\quad - \int_L^R (w\eta)_d \, dx + \int_0^d B_L \, dz \end{aligned}$$

$$\text{Let } \eta \approx \frac{\partial u}{\partial z},$$

$$0 = \int_0^d u_L \, du_L - \int_0^d u_R \, du_R - \int_L^R (w\eta)_d \, dx + \int_0^H B_L \, dz$$

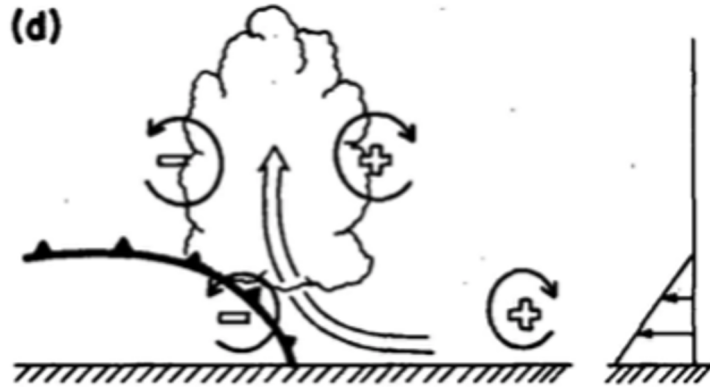
$$= \left( \overset{=0}{\frac{u_{L,d}^2}{2}} - \overset{=0}{\frac{u_{L,0}^2}{2}} \right) - \left( \frac{u_{R,d}^2}{2} - \frac{u_{R,0}^2}{2} \right) - \int_L^R (w\eta)_d \, dx$$

$$+ \int_0^H B_L \, dz, \quad \text{x = R no vertical shear}$$

$$u_{L,d}^2 = -2 \int_0^H B_L \, dz \equiv c^2$$

# Interpretation

## c. The cold pool with and without low-level shear



$$u_{L,d}^2 = -2 \int_0^H B_L dz \equiv c^2$$

$$\begin{aligned}
 0 &= \int_0^d u_L du_L - \int_0^d u_R du_R - \int_L^R (w\eta)_d dx + \int_0^H B_L dz \\
 &= \left( \frac{u_{L,d}^2}{2} - \frac{u_{L,0}^2}{2} \right) - \left( \frac{u_{R,d}^2}{2} - \frac{u_{R,0}^2}{2} \right) - \int_L^R (w\eta)_d dx \\
 &\quad + \int_0^H B_L dz,
 \end{aligned}$$

set  $u_{L,d}$ ,  $u_{R,d}$  and  $\int_L^R (w\eta)_d dx = 0$

$\Delta u = c$ , where  $\Delta u \equiv u_{R,d} - u_{R,0} = -u_{R,0}$



# Interpretation

## c. The cold pool with and without low-level shear

$$\Delta u = c, \quad \text{where} \quad \Delta u \equiv u_{R,d} - u_{R,0} = -u_{R,0}$$

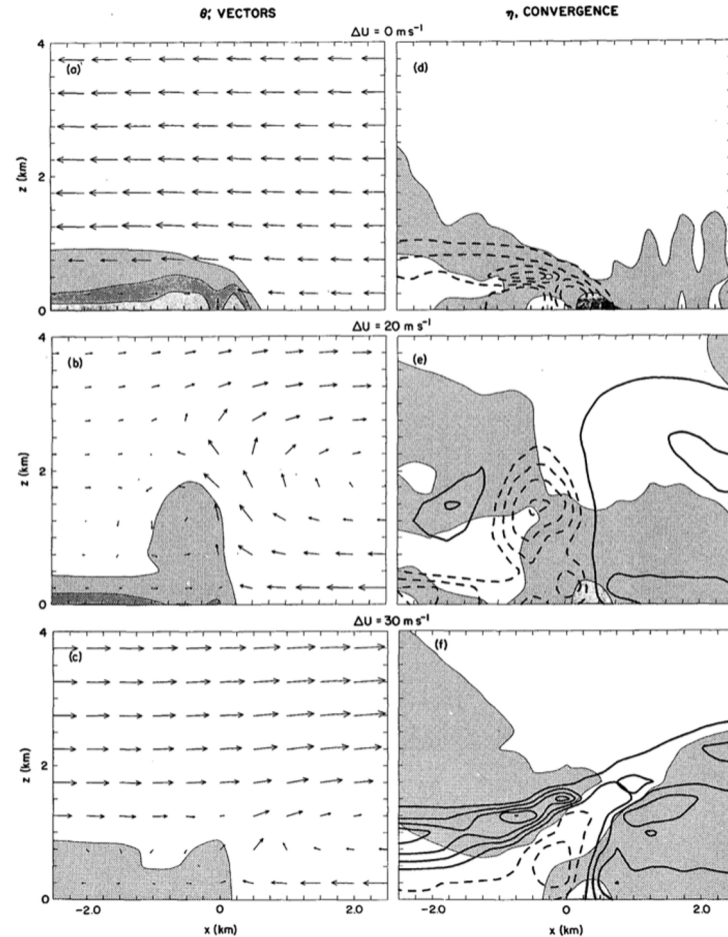
(a)-(c)

- shaded regions: negative  $\theta$  perturbation (2K intervals)

(d)-(f)

- shaded regions: convergence

- contour:  $\eta$



# Interpretation

32

## d. Speed of the cold pool

$$\begin{aligned} 0 &= \int_0^d u_L du_L - \int_0^d u_R du_R - \int_L^R (w\eta)_d dx + \int_0^H B_L dz \\ &= \left( \frac{u_{L,d}^2}{2} - \frac{u_{L,0}^2}{2} \right) - \left( \frac{u_{R,d}^2}{2} - \frac{u_{R,0}^2}{2} \right) - \int_L^R (w\eta)_d dx \\ &\quad + \int_0^H B_L dz, \end{aligned}$$

$$u_{L,d} = \epsilon u_{R,d}$$

$$\epsilon^2 u_{R,d} = \Delta u - [(1 - \epsilon^2)\Delta u^2 + \epsilon^2(c^2 + u_{L,0}^2)]^{1/2}$$

$$U_c = \begin{cases} U_{R,0} + c \left[ 1 - \frac{\beta}{2} \left( \frac{\Delta u}{c} - 1 \right)^2 \right], & \Delta u < c \\ U_{R,0} + c, & \Delta u \approx c \\ U_{R,0} + \Delta u, & \Delta u > c \end{cases}$$

$$\beta \equiv \epsilon^2 - 1$$

$$\Delta u \leq U_c - U_{R,0} \leq \max(c, \Delta u)$$

# 5

## SUMMARY

# Summary

- two basic types of simulated long-lived squall lines:
  - lines of more-or-less ordinary cells that continually grow and decay
  - lines of nearly steady supercells
- lines composed of supercell-like circulations occurring far less frequently
- sustaining a long-lived line of time-dependent cells:
  - an optimal value of wind-shear magnitude in comparison to the depth and coldness of the outflow
- the system ultimately changing to the weaker state
- a *ipso facto* steady squall line composed of supercells
- evaluating whether atmospheric condition favor the long-lived squall lines by equations

**Thanks for listening!**