

Chapter 3 Surface Integral

Assume the surface is defined by

$$\left. \begin{aligned} x &= f(u, v) \\ y &= g(u, v) \\ z &= h(u, v) \end{aligned} \right\}, \text{ or, } z = \bar{f}(x, y)$$

then let us find the surface integral $\iint_S G(x, y, z) dA = ?$

often $G(x, y, z)$ can be written as $G(x, y, z) = \mathbf{F}(x, y, z) \cdot \hat{\mathbf{n}}$

where $\hat{\mathbf{n}}$ is the normal vector of surface S ,
and \mathbf{F} is a vector function

Example: the gravitational force field is defined as

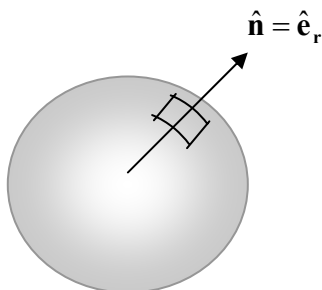
$$\mathbf{F}(r, \lambda, \phi) = -\frac{GMm}{r^2} \hat{\mathbf{e}}_r$$

The fluid velocity is defined as $\mathbf{V}(x, y, z)$

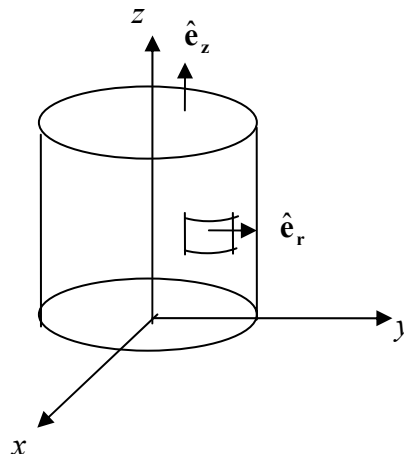
Note: $\hat{\mathbf{n}} \equiv$ unit normal vector to surface S
i.e., length = 1, and “ \perp ”(normal) to surface S

Ex:

Sphere



Cylinder



Find $\iint \mathbf{F}(x,y,z) \cdot \hat{\mathbf{n}} \, dA = ?$

\Rightarrow Get two tangent vector $\mathbf{b}_u, \mathbf{b}_v$


$$\begin{cases} \mathbf{b}_u = \frac{\partial f}{\partial u} \hat{\mathbf{i}} + \frac{\partial g}{\partial u} \hat{\mathbf{j}} + \frac{\partial h}{\partial u} \hat{\mathbf{k}} \\ \mathbf{b}_v = \frac{\partial f}{\partial v} \hat{\mathbf{i}} + \frac{\partial g}{\partial v} \hat{\mathbf{j}} + \frac{\partial h}{\partial v} \hat{\mathbf{k}} \end{cases}$$

then by using \mathbf{b}_u and \mathbf{b}_v , we can find normal vector $\hat{\mathbf{n}}$ as

$$\begin{aligned} \mathbf{N} &= \mathbf{b}_u \times \mathbf{b}_v \\ \therefore \mathbf{n} &= \pm \frac{\mathbf{b}_u \times \mathbf{b}_v}{|\mathbf{b}_u \times \mathbf{b}_v|} \end{aligned}$$

(“ \perp ” depends on the problem, but $\hat{\mathbf{n}}$ always directs outward from the surface)

$$\begin{aligned} dA &= |\mathbf{b}_u \times \mathbf{b}_v| \, du \, dv \\ \Rightarrow \iint_S \mathbf{F}(x,y,z) \cdot \hat{\mathbf{n}} \, dA \\ &= \iint_S \mathbf{F}(f(u,v), g(u,v), h(u,v)) \cdot (\mathbf{b}_u \times \mathbf{b}_v) \, du \, dv \quad (3.1) \end{aligned}$$


 need to compute each component

Ex: For $z = \bar{f}(x, y)$,

$$\begin{aligned} \Rightarrow \begin{cases} x = x \\ y = y \\ z = \bar{f}(x, y) \end{cases} & \quad \begin{aligned} \mathbf{b}_x &= \hat{\mathbf{i}} + \frac{\partial \bar{f}}{\partial x} \hat{\mathbf{k}} \\ \mathbf{b}_y &= \hat{\mathbf{j}} + \frac{\partial \bar{f}}{\partial y} \hat{\mathbf{k}} \end{aligned} \end{aligned}$$

$$\Rightarrow \mathbf{b}_x \times \mathbf{b}_y = \mathbf{N} = \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ 1 & 0 & \frac{\partial \bar{f}}{\partial x} \\ 0 & 1 & \frac{\partial \bar{f}}{\partial y} \end{vmatrix} = -\frac{\partial \bar{f}}{\partial x} \hat{\mathbf{i}} - \frac{\partial \bar{f}}{\partial y} \hat{\mathbf{j}} + \hat{\mathbf{k}}$$

$$\Rightarrow \hat{\mathbf{n}} = \frac{\mathbf{N}}{\left[\left(\frac{\partial \bar{f}}{\partial x} \right)^2 + \left(\frac{\partial \bar{f}}{\partial y} \right)^2 + 1 \right]^{1/2}} = \frac{-\frac{\partial \bar{f}}{\partial x} \hat{\mathbf{i}} - \frac{\partial \bar{f}}{\partial y} \hat{\mathbf{j}} + \hat{\mathbf{k}}}{\left[\left(\frac{\partial \bar{f}}{\partial x} \right)^2 + \left(\frac{\partial \bar{f}}{\partial y} \right)^2 + 1 \right]^{1/2}}$$

$$\therefore \iint_S \mathbf{F}(x,y,z) \cdot \hat{\mathbf{n}} \, dA = \iint_{S(x,y)} \mathbf{F}(x,y,\bar{f}(x,y)) \cdot \left(-\frac{\partial \bar{f}}{\partial x}, -\frac{\partial \bar{f}}{\partial y}, 1 \right) \, dx dy$$

Note: the physical meaning of $\mathbf{F} \cdot \hat{\mathbf{n}} \equiv$ the flux of \mathbf{F} through surface S

Ex: Consider a vector field $\mathbf{F} = \hat{\mathbf{i}} + xy \hat{\mathbf{j}}$

$$\text{a surface } S: \left. \begin{array}{l} x = u + v \\ y = u - v \\ z = u^2 \end{array} \right\} \begin{array}{l} 0 \leq u \leq 1 \\ 0 \leq v \leq 1 \end{array}$$

Find the surface integral $\iint_S \mathbf{F} \cdot \hat{\mathbf{n}} \, dA = ?$

$$\text{Sol: } \because \frac{x+y}{2} = u \quad \therefore z = \frac{(x+y)^2}{4}$$

$$\mathbf{b}_u = \frac{\partial \mathbf{f}}{\partial u} \hat{\mathbf{i}} + \frac{\partial \mathbf{g}}{\partial u} \hat{\mathbf{j}} + \frac{\partial \mathbf{h}}{\partial u} \hat{\mathbf{k}} = \hat{\mathbf{i}} + \hat{\mathbf{j}} + 2u \hat{\mathbf{k}}$$

$$\mathbf{b}_v = \frac{\partial \mathbf{f}}{\partial v} \hat{\mathbf{i}} + \frac{\partial \mathbf{g}}{\partial v} \hat{\mathbf{j}} + \frac{\partial \mathbf{h}}{\partial v} \hat{\mathbf{k}} = \hat{\mathbf{i}} - \hat{\mathbf{j}}$$

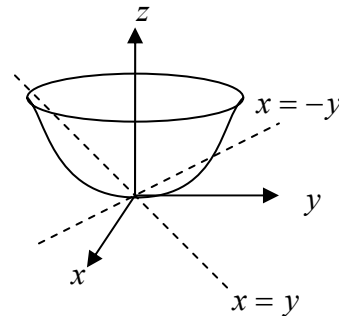
$$\therefore \iint_S \mathbf{F} \cdot \hat{\mathbf{n}} \, dA = \iint_{S(u,v)} \mathbf{F} \cdot (\mathbf{b}_u \times \mathbf{b}_v) \, dudv$$

$$= \iint \left[\hat{\mathbf{i}} + (u^2 - v^2) \hat{\mathbf{j}} \right] \cdot [\mathbf{b}_u \times \mathbf{b}_v] \, dudv$$

$$= \int_0^1 \int_0^1 [1, u^2 - v^2, 0] \cdot [2u, 2u, -2] \, dudv$$

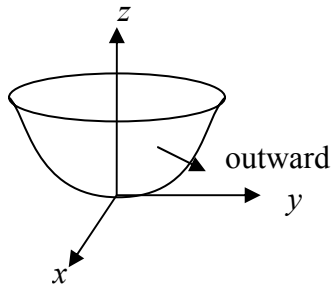
$$= \int_0^1 \int_0^1 [2u + 2u(u^2 - v^2)] \, dudv = \int_0^1 \int_0^1 [2u + 2u^3 - 2uv^2] \, dudv$$

$$= \int_0^1 \left[u^2 + \frac{1}{2} u^4 - u^2 v^2 \right]_{u=0}^{u=1} dv = \int_0^1 \left[1 + \frac{1}{2} - v^2 \right] dv = \frac{3}{2} - \frac{1}{3} = \frac{7}{6}$$



$$\left(\mathbf{b}_u \times \mathbf{b}_v = \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ 1 & 1 & 2u \\ 1 & -1 & 0 \end{vmatrix} = +2u\hat{\mathbf{i}} + 2u\hat{\mathbf{j}} + (-2u)\hat{\mathbf{k}} \right)$$

Ex:



$$x^2 + y^2 \leq 1, \quad z = \frac{x^2 + y^2}{4}$$

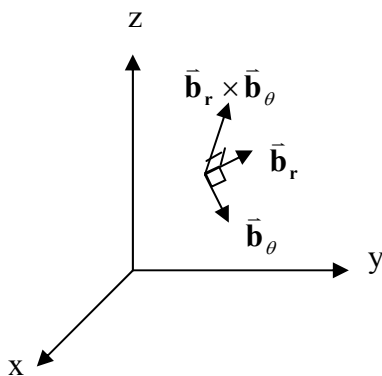
$$\text{For } \mathbf{F} = \hat{\mathbf{i}} + xy\hat{\mathbf{j}}$$

$$\text{Find } \iint \mathbf{F} \cdot \hat{\mathbf{n}} \, dA = ?$$

Sol: Transform the Cartesian coordinate into the cylindrical coordinate

$$0 \leq r \leq 1 \quad 0 \leq \theta \leq 2\pi \Rightarrow \begin{cases} x = r \cos \theta = f(r, \theta) \\ y = r \sin \theta = g(r, \theta) \\ z = \frac{x^2 + y^2}{4} = \frac{r^2}{4} = h(r, \theta) \end{cases}$$

$$\Rightarrow \begin{aligned} \mathbf{b}_r &= \frac{\partial f}{\partial r} \hat{\mathbf{i}} + \frac{\partial g}{\partial r} \hat{\mathbf{j}} + \frac{\partial h}{\partial r} \hat{\mathbf{k}} = \cos \theta \hat{\mathbf{i}} + \sin \theta \hat{\mathbf{j}} + \frac{r}{2} \hat{\mathbf{k}} \\ \mathbf{b}_\theta &= \frac{\partial f}{\partial \theta} \hat{\mathbf{i}} + \frac{\partial g}{\partial \theta} \hat{\mathbf{j}} + \frac{\partial h}{\partial \theta} \hat{\mathbf{k}} = -r \sin \theta \hat{\mathbf{i}} + r \cos \theta \hat{\mathbf{j}} + 0 \hat{\mathbf{k}} \end{aligned}$$



$$\text{Thus, Find } \iint \mathbf{F} \cdot \hat{\mathbf{n}} \, dA = \iint \mathbf{F} \cdot (\mathbf{b}_\theta \times \mathbf{b}_r) \, dr \, d\theta$$

Note that

$$\begin{aligned} \mathbf{F} \cdot (\mathbf{b}_\theta \times \mathbf{b}_r) &= \begin{vmatrix} 1 & r^2 \cos \theta \sin \theta & 0 \\ -r \sin \theta & r \cos \theta & 0 \\ \cos \theta & \sin \theta & \frac{r}{2} \end{vmatrix} \\ &= \frac{r^2}{2} \cos \theta - r^2 \cos \theta \sin \theta \left(-\frac{r^2}{2} \sin \theta \right) \\ &= \frac{r^2}{2} \cos \theta + \frac{r^4}{2} \cos \theta \sin^2 \theta \end{aligned}$$

$$\begin{aligned}
\Rightarrow \iint \mathbf{F} \cdot \hat{\mathbf{n}} dA &= \int_0^{2\pi} \int_0^1 \left[\frac{r^2}{2} \cos \theta + \frac{r^4}{2} \cos \theta \sin^2 \theta \right] dr d\theta \\
&= \int_0^{2\pi} \left[\frac{1}{6} \cos \theta + \frac{1}{10} \cos \theta \sin^2 \theta \right] d\theta \\
&= \frac{1}{6} \int_0^{2\pi} \cos \theta d\theta + \frac{1}{10} \int_0^{2\pi} \sin^2 \theta d\theta \\
&= \frac{1}{6} \cdot 0 + \frac{1}{30} (0 - 0) = 0
\end{aligned}$$

● Divergence of Vector fields

Recall the Gauss's Law of Electrostatics:

$$\iint_S \mathbf{E} \cdot \hat{\mathbf{n}} dA = 4\pi q_o \quad (3.2)$$

where S is any closed surface,

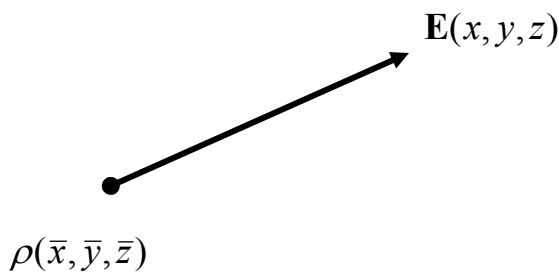
\mathbf{E} is the electrical field,

$\hat{\mathbf{n}}$ is the normal vector of surface S,

q_o is the total electrical charge contained by surface S.

Another way to find the electrical field \mathbf{E} :

$$\mathbf{E}(x, y, z) = \iiint \frac{\rho(\bar{x}, \bar{y}, \bar{z}) \hat{u}(r - \bar{r})}{(x - \bar{x})^2 + (y - \bar{y})^2 + (z - \bar{z})^2} d\bar{x} d\bar{y} d\bar{z}$$



where $\rho(\bar{x}, \bar{y}, \bar{z})$ is the charge density

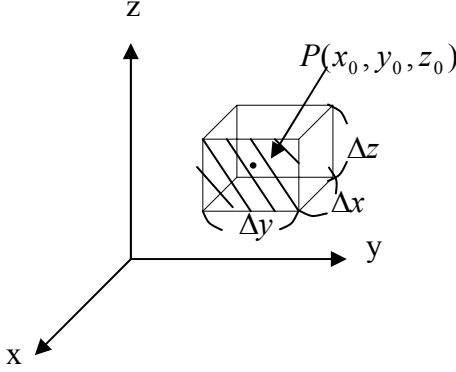
Consider \mathbf{E} changes smoothly and ρ also changes smoothly

$$\text{Thus } \iint_S \mathbf{E} \cdot \hat{\mathbf{n}} dA = 4\pi q_o \approx 4\pi \rho(x, y, z) \Delta V$$

$$\Rightarrow \lim_{\Delta V \rightarrow 0} \frac{1}{\Delta V} \iint_{\Delta S} \mathbf{E} \cdot \hat{\mathbf{n}} dA = 4\pi \rho(x, y, z) = \text{div} \mathbf{E} = \nabla \cdot \mathbf{E}$$

Thus

$$\boxed{\lim_{\Delta V \rightarrow 0} \frac{1}{\Delta V} \iint_{\Delta S} \mathbf{E} \cdot \hat{\mathbf{n}} dA = \nabla \cdot \mathbf{E}} \quad (3.3)$$



$\Delta V = \Delta x \Delta y \Delta z$, $P(x_0, y_0, z_0)$: center of cubic

Consider $\mathbf{E} = E_x \hat{\mathbf{i}} + E_y \hat{\mathbf{j}} + E_z \hat{\mathbf{k}}$

for the front surface $\Rightarrow \hat{\mathbf{n}} = \hat{\mathbf{i}}$

for the back surface $\Rightarrow \hat{\mathbf{n}} = -\hat{\mathbf{i}}$

Then at the front surface,

$$\mathbf{E} \cdot \hat{\mathbf{n}} = \mathbf{E} \left(x_0 + \frac{\Delta x}{2}, y_0, z_0 \right) \cdot (\hat{\mathbf{i}}) = E_x \left(x_0 + \frac{\Delta x}{2} \right)$$

$$\Delta A = \Delta y \Delta z$$

$$\therefore \iint_{\text{FRONT}} \mathbf{E} \cdot \hat{\mathbf{n}} dA = E_x \left(x_0 + \frac{\Delta x}{2} \right) \Delta y \Delta z$$

$$\text{Similarly, } \iint_{\text{BACK}} \mathbf{E} \cdot \hat{\mathbf{n}} dA = -E_x \left(x_0 - \frac{\Delta x}{2} \right) \Delta y \Delta z$$

$$\Rightarrow \iint_{\text{FRONT}} \mathbf{E} \cdot \hat{\mathbf{n}} dA + \iint_{\text{BACK}} \mathbf{E} \cdot \hat{\mathbf{n}} dA = \left[E_x \left(x_0 + \frac{\Delta x}{2} \right) - E_x \left(x_0 - \frac{\Delta x}{2} \right) \right] \Delta y \Delta z \quad (\text{A})$$

$$\frac{(\text{A})}{\Delta V} \Rightarrow \frac{1}{\Delta V} \left\{ \iint_{\text{FRONT}} \mathbf{E} \cdot \hat{\mathbf{n}} dA + \iint_{\text{BACK}} \mathbf{E} \cdot \hat{\mathbf{n}} dA \right\} = \frac{1}{\Delta x} \left[E_x \left(x_0 + \frac{\Delta x}{2} \right) - E_x \left(x_0 - \frac{\Delta x}{2} \right) \right]$$

$$(\Delta V = \Delta x \Delta y \Delta z) \left(\begin{array}{l} \text{further assume} \\ \Delta V \rightarrow 0 \text{ i.e., } \Delta x \rightarrow 0, \\ \Delta y \rightarrow 0, \Delta z \rightarrow 0 \end{array} \right) = \frac{\partial E_x}{\partial x} (x_0, y_0, z_0)$$

Similarly for other two surface (LEFT & RIGHT, UP & DOWN)

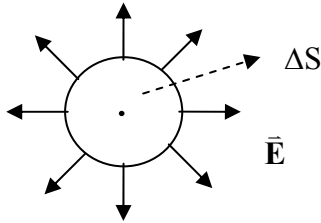
$$\Rightarrow \nabla \cdot \mathbf{E} = \lim_{\Delta V \rightarrow 0} \frac{1}{\Delta V} \iint_{\Delta S} \mathbf{E} \cdot \hat{\mathbf{n}} dA = \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} \quad (3.4)$$

where $\mathbf{E} = E_x \hat{\mathbf{i}} + E_y \hat{\mathbf{j}} + E_z \hat{\mathbf{k}}$

∴ Gauss's Law can be written as

$$\boxed{\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 4\pi\rho(x, y, z)} \quad (3.5)$$

Physical meaning of $\nabla \cdot \mathbf{E}$:



$$\nabla \cdot \mathbf{E} = \lim_{\Delta V \rightarrow 0} \frac{1}{\Delta V} \iint_{\Delta S} \mathbf{E} \cdot \hat{\mathbf{n}} dA$$

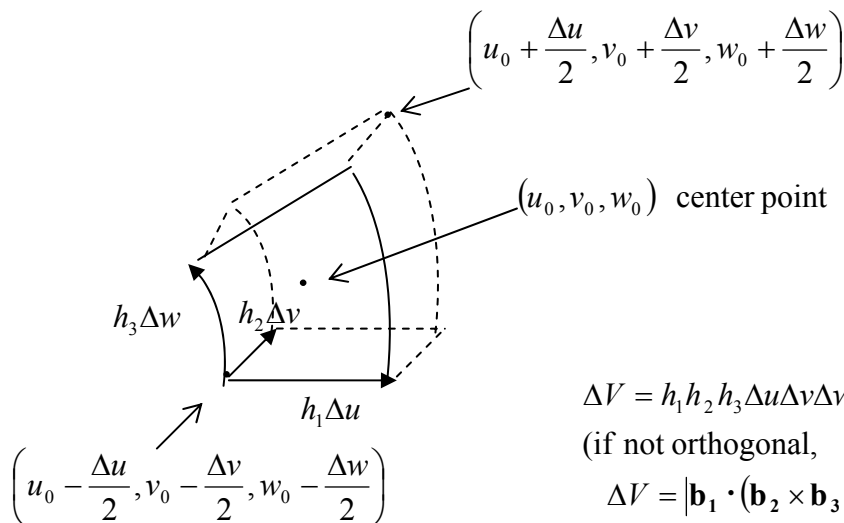
Divergence of vector field \mathbf{E} (per unit volume)

Divergence of \mathbf{F} in Cartesian coordinate:

$$\nabla \cdot \mathbf{F} = \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} + \frac{\partial F_z}{\partial z}, \quad \text{where } \mathbf{F} = F_x \hat{\mathbf{i}} + F_y \hat{\mathbf{j}} + F_z \hat{\mathbf{k}}$$

For a general orthogonal coordinate system (u, v, w)

$$\begin{cases} x = f(u, v, w) \\ y = g(u, v, w) \\ z = h(u, v, w) \end{cases} \quad \mathbf{b}_1, \mathbf{b}_2, \mathbf{b}_3 \text{ with length } h_1, h_2, h_3$$



$$\Delta V = h_1 h_2 h_3 \Delta u \Delta v \Delta w \quad (\because \text{orthogonal})$$

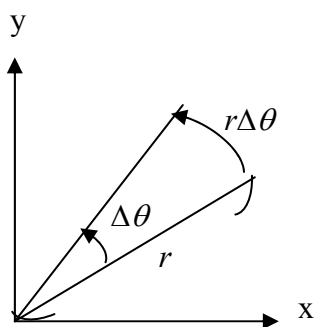
(if not orthogonal,

$$\Delta V = |\mathbf{b}_1 \cdot (\mathbf{b}_2 \times \mathbf{b}_3)| du dv dw$$

By definition,

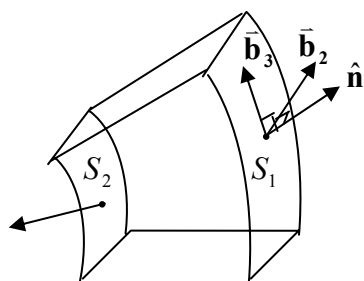
$$\begin{aligned}\nabla \cdot \mathbf{F} &= \lim_{\Delta V \rightarrow 0} \frac{1}{\Delta V} \iint_S \mathbf{F} \cdot \hat{\mathbf{n}} \, dA \\ &= \lim_{\Delta V \rightarrow 0} \frac{1}{h_1 h_2 h_3 \Delta u \Delta v \Delta w} \iint_S \mathbf{F} \cdot \hat{\mathbf{n}} \, dA\end{aligned}$$

Ex: Consider the polar coordinate (r, θ)



$$h_1 \Delta u = \Delta r \quad (h_1 = 1)$$

$$h_2 \Delta v = r \Delta \theta \quad (h_2 = r)$$



Consider these two surfaces

$$S_1 : "u = u_0 + \frac{\Delta u}{2}" \text{ constant surface}$$

$$S_2 : "u = u_0 - \frac{\Delta u}{2}" \text{ constant surface}$$

$$\hat{\mathbf{n}} = \frac{\mathbf{b}_2 \times \mathbf{b}_3}{|\mathbf{b}_2 \times \mathbf{b}_3|}, \quad \Delta A = |\mathbf{b}_2 \times \mathbf{b}_3| \Delta v \Delta w$$

$$\begin{aligned}\text{Then } \iint_S \mathbf{F} \cdot \hat{\mathbf{n}} \, dA &= F\left(u_0 + \frac{\Delta u}{2}, v_0, w_0\right) \cdot \frac{\mathbf{b}_2 \times \mathbf{b}_3}{|\mathbf{b}_2 \times \mathbf{b}_3|} |\mathbf{b}_2 \times \mathbf{b}_3| \Delta v \Delta w \\ &= F\left(u_0 + \frac{\Delta u}{2}, v_0, w_0\right) \cdot (\mathbf{b}_2 \times \mathbf{b}_3) \Delta v \Delta w\end{aligned}$$

Decompose \mathbf{F} into three components:

$$\mathbf{F} = F_u \hat{\mathbf{e}}_u + F_v \hat{\mathbf{e}}_v + F_w \hat{\mathbf{e}}_w$$

Along surface S_1 :

$$\hat{\mathbf{n}} = \hat{\mathbf{e}}_u, \quad \mathbf{b}_2 \times \mathbf{b}_3 = h_2 h_3 \hat{\mathbf{e}}_u$$

$$\therefore \iint_{S_1} \mathbf{F} \cdot \hat{\mathbf{n}} \, dA = F_u \left(u_0 + \frac{\Delta u}{2}, v_0, w_0 \right) h_2 \left(u_0 + \frac{\Delta u}{2}, v_0, w_0 \right) h_3 \left(u_0 + \frac{\Delta u}{2}, v_0, w_0 \right) \Delta v \Delta w$$

$$\text{Similarly } \iint_{S_2} \mathbf{F} \cdot \hat{\mathbf{n}} \, dA = -F_u \left(u_0 - \frac{\Delta u}{2}, v_0, w_0 \right) h_2 \left(u_0 - \frac{\Delta u}{2}, v_0, w_0 \right) h_3 \left(u_0 - \frac{\Delta u}{2}, v_0, w_0 \right) \Delta v \Delta w$$

$$\begin{aligned} \text{Thus, } \lim_{\Delta V \rightarrow 0} \frac{1}{\Delta V} & \left\{ \iint_{S_1} \mathbf{F} \cdot \hat{\mathbf{n}} \, dA + \iint_{S_2} \mathbf{F} \cdot \hat{\mathbf{n}} \, dA \right\} \\ &= \lim_{\Delta V \rightarrow 0} \frac{\Delta v \Delta w}{h_1 h_2 h_3 \Delta u \Delta v \Delta w} \left\{ F_u \left(u_0 + \frac{\Delta u}{2}, v_0, w_0 \right) h_2 \left(u_0 + \frac{\Delta u}{2}, v_0, w_0 \right) h_3 \left(u_0 + \frac{\Delta u}{2}, v_0, w_0 \right) \right. \\ & \quad \left. - F_u \left(u_0 - \frac{\Delta u}{2}, v_0, w_0 \right) h_2 \left(u_0 - \frac{\Delta u}{2}, v_0, w_0 \right) h_3 \left(u_0 - \frac{\Delta u}{2}, v_0, w_0 \right) \right\} \\ &= \frac{1}{h_1 h_2 h_3} \frac{\partial}{\partial u} (F_u h_2 h_3) \end{aligned}$$

Similar derivations are for other four surfaces (UP & DOWN, FRONT & BACK)

Thus the final expression is

$$\nabla \cdot \mathbf{F} = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u} (F_u h_2 h_3) + \frac{\partial}{\partial v} (F_v h_3 h_1) + \frac{\partial}{\partial w} (F_w h_1 h_2) \right] \quad (3.6)$$

Ex: For the Cylindrical coordinate (r, θ, z) :

$$h_1 = 1 \Leftrightarrow r \text{ - component}$$

$$h_2 = r \Leftrightarrow \theta \text{ - component}$$

$$h_3 = 1 \Leftrightarrow z \text{ - component}$$

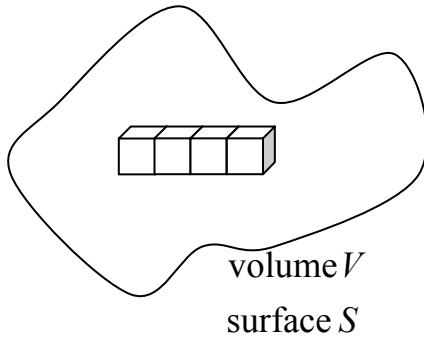
$$\begin{aligned} \therefore \nabla \cdot \mathbf{F} &= \frac{1}{r} \left[\frac{\partial}{\partial r} (F_r \cdot r) + \frac{\partial}{\partial \theta} (F_\theta) + \frac{\partial}{\partial z} (F_z \cdot r) \right] \\ &= \frac{1}{r} \left[r \frac{\partial F_r}{\partial r} + F_r \right] + \frac{1}{r} \frac{\partial F_\theta}{\partial \theta} + \frac{\partial F_z}{\partial z} \\ &= \frac{\partial F_r}{\partial r} + \frac{1}{r} F_r + \frac{1}{r} \frac{\partial F_\theta}{\partial \theta} + \frac{\partial F_z}{\partial z} \end{aligned}$$

Divergence Theorem (or The Gauss Theorem):

$$\iint_S \mathbf{F} \cdot \hat{\mathbf{n}} \, dA = \iiint_V \nabla \cdot \mathbf{F} \, dV \quad (3.7)$$

where S is a surface

V is the volume contained by the closed surface S



$$\begin{aligned} \iint_S \mathbf{F} \cdot \hat{\mathbf{n}} \, dA &= \sum_{\substack{\text{small} \\ \text{volume}}} \iint_{S_i} \mathbf{F} \cdot \hat{\mathbf{n}} \, dA \\ &= \lim_{\Delta V_i \rightarrow 0} \sum_{\substack{\text{small} \\ \text{volume}}} \Delta V_i \left(\frac{1}{\Delta V_i} \iint_{S_i} \mathbf{F} \cdot \hat{\mathbf{n}} \, dA \right) \\ &= \lim_{\Delta V_i \rightarrow 0} \sum_{\substack{\text{small} \\ \text{volume}}} \Delta V_i (\nabla \cdot \mathbf{F}) \\ &= \iiint_V \nabla \cdot \mathbf{F} \, dV \end{aligned}$$

Application of divergence theorem \Rightarrow mass conservation

density = $\rho(x, y, z, t)$

Let velocity = $\vec{\mathbf{V}}(x, y, z, t)$

$M(t)$ = total mass in the fixed volume

$$\text{then } \frac{dM(t)}{dt} = \frac{d}{dt} \iiint_V \rho(x, y, z, t) \, dV = \iiint_V \frac{\partial \rho}{\partial t} \, dV$$

\uparrow
 fixed volume V

In other words,

$$\begin{aligned} \frac{dM(t)}{dt} &= - \iint_S \rho \mathbf{V} \cdot \hat{\mathbf{n}} \, dA && \text{(flux outward of } \rho \mathbf{V}) \\ &= - \iiint_V \nabla \cdot (\rho \mathbf{V}) \, dV && \text{(using the divergence theorem)} \end{aligned}$$

$$\therefore \iiint_V \frac{\partial \rho}{\partial t} \, dV = - \iiint_V \nabla \cdot (\rho \mathbf{V}) \, dV$$

$$\text{Thus for any volume } V \Rightarrow \frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \mathbf{V}) \quad (3.8) \quad \text{Continuity Equation}$$

Note that for incompressible fluid, $\rho = \text{constant}$, i.e., $\frac{\partial \rho}{\partial t} = 0$

So the continuity equation is reduced to

$$0 = \nabla \cdot (\rho \mathbf{V}) = \rho \nabla \cdot \mathbf{V} \Rightarrow \nabla \cdot \mathbf{V} = 0 \quad (3.9)$$

Continuity Equation for Incompressible fluid